

# Electroweak Sudakov Corrections to Scattering Processes

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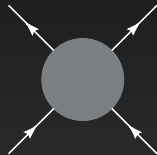
Based on:

J.C., F. Golf, R. Kelley, A. Manohar, PRL 100 (2008) 021802

J.C., F. Golf, R. Kelley, A. Manohar, PRD 77 (2008) 053004

J.C., R. Kelley, A. Manohar in preparation

- $2 \rightarrow 2$  scattering – one of the dominant high-energy processes at LHC.
- The EFT operators become 4 particle operators.
- Shown that we can use previous results, by summing over pairs of particles, with appropriate group theory factors.
- Can be used for:
  - ▶ quark production:  $q\bar{q} \rightarrow q\bar{q}$  or  $q\bar{q} \rightarrow t\bar{t}$
  - ▶ squark production:  $q\bar{q} \rightarrow \tilde{t}\tilde{t}^*$
  - ▶ gauge boson production:  $q\bar{q} \rightarrow ZZ, WW, \dots$



$$\alpha = \alpha_2 = \frac{\alpha_{em}}{\sin^2 \theta_W}$$

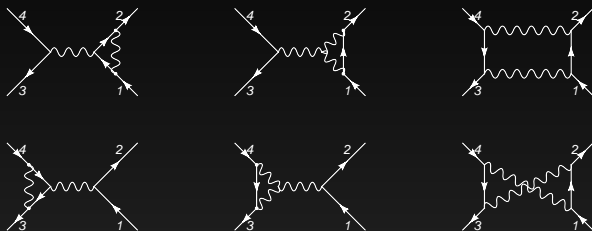
$$C(Q, \mu) \xrightarrow[\text{SCET w/ } M=0]{\text{Full Theory}} \mu = Q$$



$$D(M, \mu) \xrightarrow[\text{SCET w/o gauge boson}]{\text{SCET w/ mass } M} \mu = M$$

# 4-particle full theory

- Full Theory:  $\mathcal{O}_{full} = \left[ \bar{\psi} \gamma^\mu \mathbf{t}^a \psi \right] \left[ \bar{\psi} \gamma^\mu \mathbf{t}^a \psi \right]$
- One-loop diagrams:



$$\mathbf{t}^a \mathbf{t}^a = C_F \mathbf{1}; \quad \mathbf{t}^a \mathbf{t}^b \mathbf{t}^a = \left( C_F - \frac{1}{2} C_A \right) \mathbf{t}^b$$

$$\mathbf{t}^a \mathbf{t}^b \otimes \mathbf{t}^a \mathbf{t}^b = C_1 \mathbf{1} \otimes \mathbf{1} + \frac{1}{4} \left( C_d - \frac{1}{2} C_A \right) \mathbf{t}^a \otimes \mathbf{t}^a$$

$$\mathbf{t}^a \mathbf{t}^b \otimes \mathbf{t}^b \mathbf{t}^a = C_1 \mathbf{1} \otimes \mathbf{1} + \frac{1}{4} \left( C_d + \frac{1}{2} C_A \right) \mathbf{t}^a \otimes \mathbf{t}^a$$

# 4-particle SCET

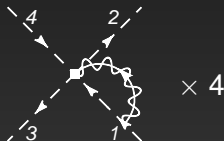
$$\begin{aligned}
 \bullet \mathcal{O}_{SCET} &= C_1 \left[ [\bar{\xi}_{n,p_2} W_n] \gamma^\mu t^a [W_{\bar{n}}^\dagger \xi_{\bar{n},p_1}] \right] \left[ [\bar{\xi}_{n,p_4} W_n] \gamma^\mu t^a [W_{\bar{n}}^\dagger \xi_{\bar{n},p_3}] \right] \\
 &\quad + C_2 \left[ [\bar{\xi}_{n,p_2} W_n] \gamma^\mu [W_{\bar{n}}^\dagger \xi_{\bar{n},p_1}] \right] \left[ [\bar{\xi}_{n,p_4} W_n] \gamma^\mu [W_{\bar{n}}^\dagger \xi_{\bar{n},p_3}] \right] \\
 &= C_1 [\mathbf{t}^a \otimes \mathbf{t}^a]_{SCET} + C_2 [\mathbf{1} \otimes \mathbf{1}]_{SCET}
 \end{aligned}$$

- One Loop : 4 Collinear diagrams + 6 Ultra-Soft diagrams + wavefunction corrections

- Ultra-soft Graphs



- Collinear Graphs



# Matrix Element in Matrix Form (in $C_1, C_2$ space)

$$R = \tilde{R}\mathbf{1} + R_S \quad \Gamma_{ij} = \mathbf{1}_{n_i} + \mathbf{1}_{n_j} + \mathbf{1}_{s_{ij}}$$

$$\tilde{R} = C_F[\Gamma_{12} - \frac{1}{2}\delta Z_1^{-1} - \frac{1}{2}\delta Z_2^{-1}] + C_F[\Gamma_{34} - \frac{1}{2}\delta Z_3^{-1} - \frac{1}{2}\delta Z_4^{-1}]$$

$$R_S = \begin{bmatrix} \frac{1}{4}C_d r_1 + \frac{1}{4}C_A r_2 & r_1 \\ C_1 r_1 & 0 \end{bmatrix}$$

$$r_1 = \Gamma_{14} + \Gamma_{23} - \Gamma_{13} - \Gamma_{24}$$

$$r_2 = \Gamma_{14} + \Gamma_{23} + \Gamma_{13} + \Gamma_{24} - \Gamma_{12} - \Gamma_{34}$$

$$\Gamma_{ij} \rightarrow S_{ij} \equiv \Gamma_{ij} - \frac{1}{2}\delta Z_i^{-1} - \frac{1}{2}\delta Z_j^{-1}$$

$$\implies \tilde{R} = C_F S_{12} + C_F S_{34}$$

$$r_1 = S_{14} + S_{23} - S_{13} - S_{24}$$

$$r_2 = S_{14} + S_{23} + S_{13} + S_{24} - S_{12} - S_{34}$$

(also true for arbitrary number of external particles)

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# SU(2) Toy Theory

anomalous dimension:

$$\gamma^{(1)} = \tilde{\gamma}^{(1)} \mathbf{1} + \gamma_S^{(1)}$$

$$\tilde{\gamma}^{(1)} = 2C_F \left( 4 \log \frac{-s}{\mu^2} - 6 \right)$$

$$\gamma_S^{(1)} = \begin{bmatrix} 2C_d \log \frac{t}{u} + 2C_A \log \frac{ut}{s^2} & 8 \log \frac{t}{u} \\ 8C_1 \log \frac{t}{u} & 0 \end{bmatrix}$$

$\tilde{\gamma}^{(1)}$  is twice the Sudakov  $\gamma$ , has a  $L_Q$

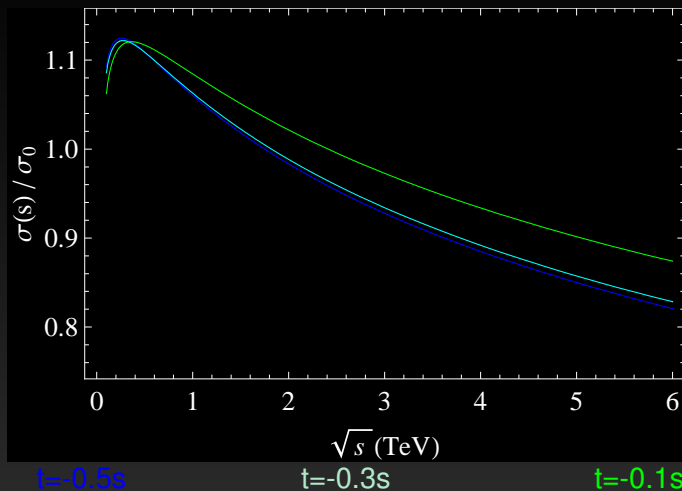
$\gamma_S^{(1)}$  has no large logarithm since  $s$ ,  $t$ , and  $u$  are all at the order of  $Q$ .

Matching Matrix at low scale:  
(same structure as anomalous dimension)

$$R^{(1)} = \tilde{R}^{(1)} \mathbf{1} + R_S^{(1)}$$

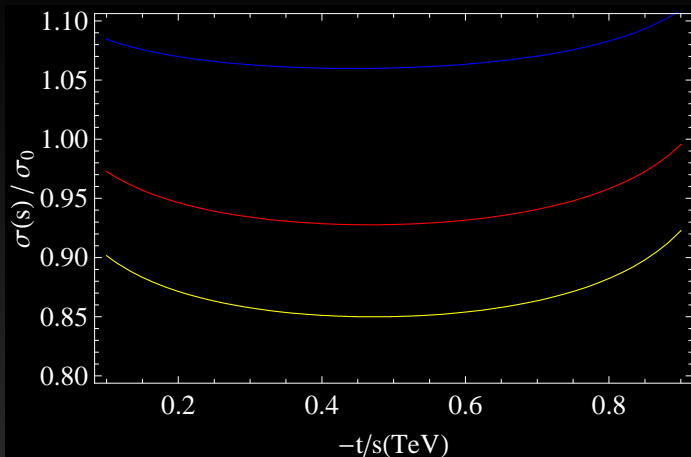
$$\tilde{R}^{(1)} = 2C_F \left( -L_M^2 + 2L_M \log \frac{-s}{\mu^2} - 2L_M + \frac{9}{2} - \frac{5\pi^2}{6} \right)$$

$$R_S^{(1)} = L_M \begin{bmatrix} C_d \log \frac{t}{u} + C_A \log \frac{ut}{s^2} & 4 \log \frac{t}{u} \\ 4C_1 \log \frac{t}{u} & 0 \end{bmatrix}$$



$\propto \frac{M^2}{Q^2}$  term neglected.

# cross-section in toy theory



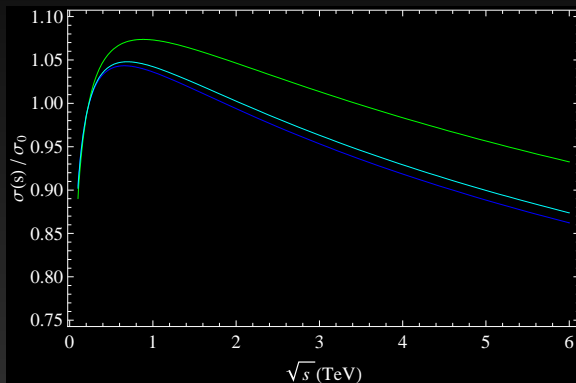
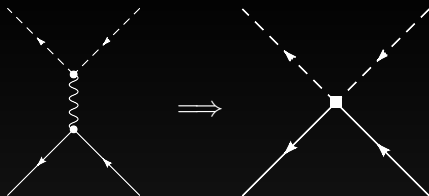
$s = 1 \text{ TeV}$

$s = 3 \text{ TeV}$

$s = 5 \text{ TeV}$

# squarks production

in SU(2) toy theory,  
taking  $m_s = 500\text{ GeV}$

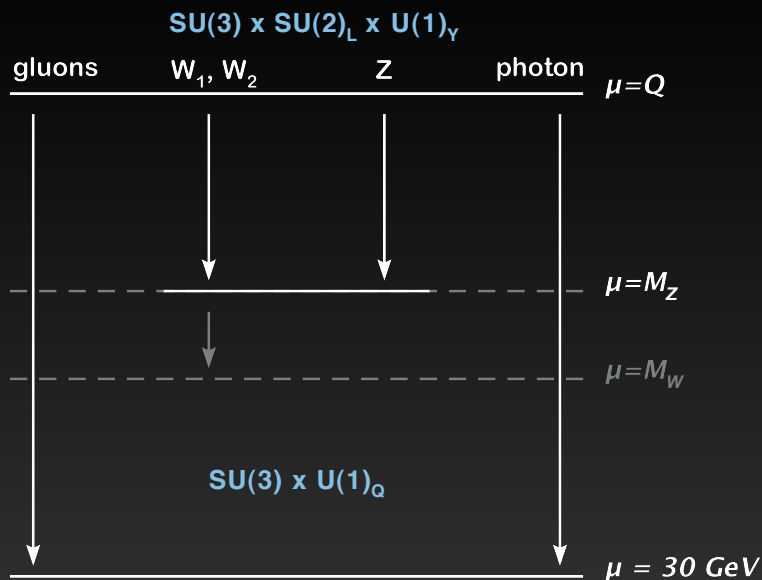


$t = -0.5\text{ s}$

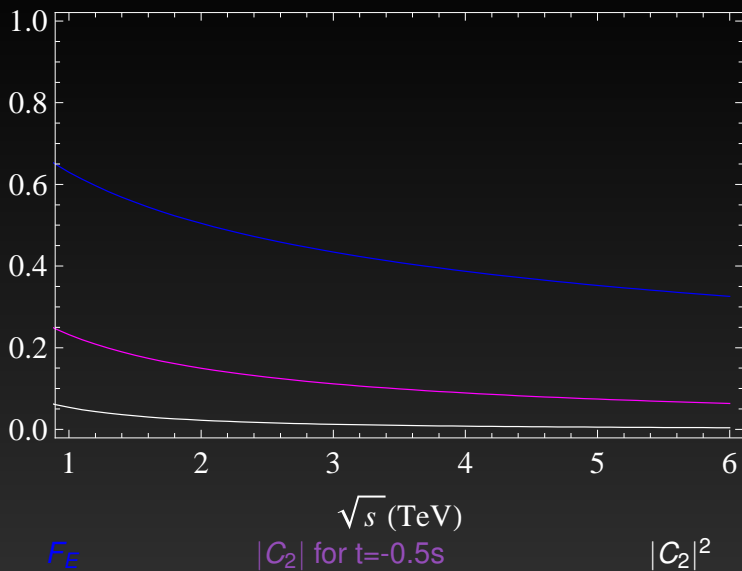
$t = -0.3\text{ s}$

$t = -0.1\text{ s}$

# Standard Model (light quarks):

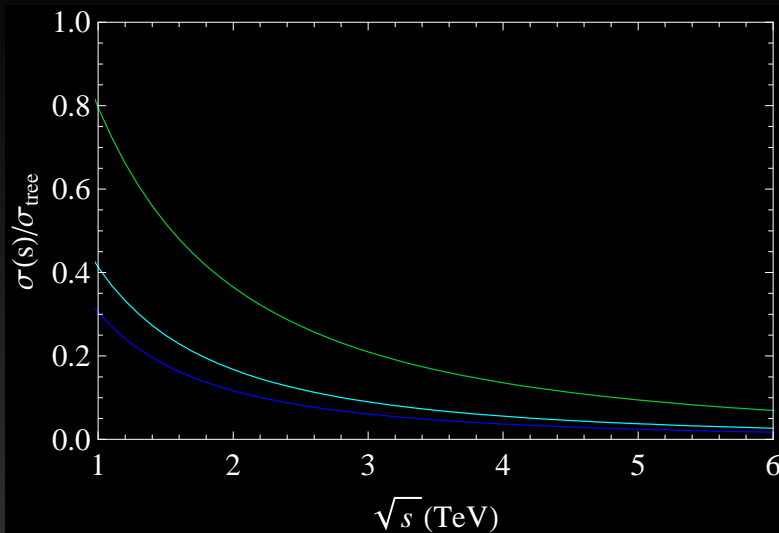


# light quarks Sudakov vs. $u\bar{u}_L \rightarrow c\bar{c}_L$



$$u\bar{u}_L \rightarrow c\bar{c}_L$$

$$\sigma = \left(\frac{2}{9}|C_1|^2 + |C_2|^2\right) \times \sigma_0$$

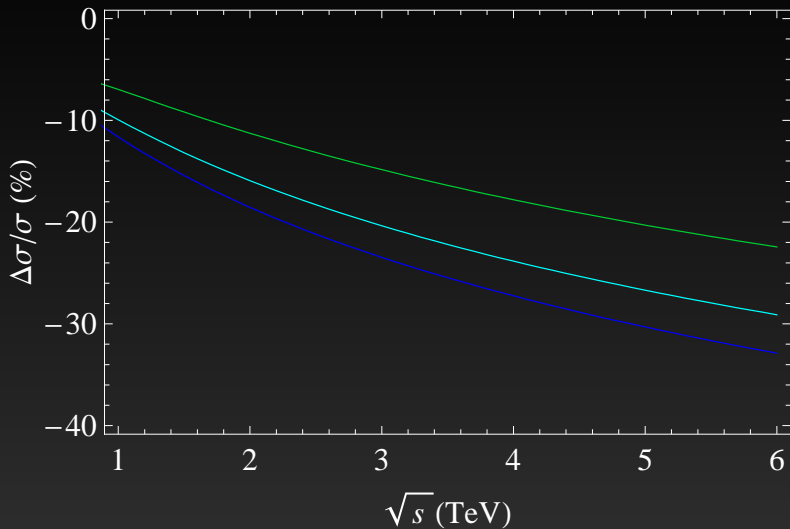


$t = -0.5s$

$t = -0.3s$

$t = -0.1s$

# $u\bar{u}_L \rightarrow c\bar{c}_L$ , contribution from EW

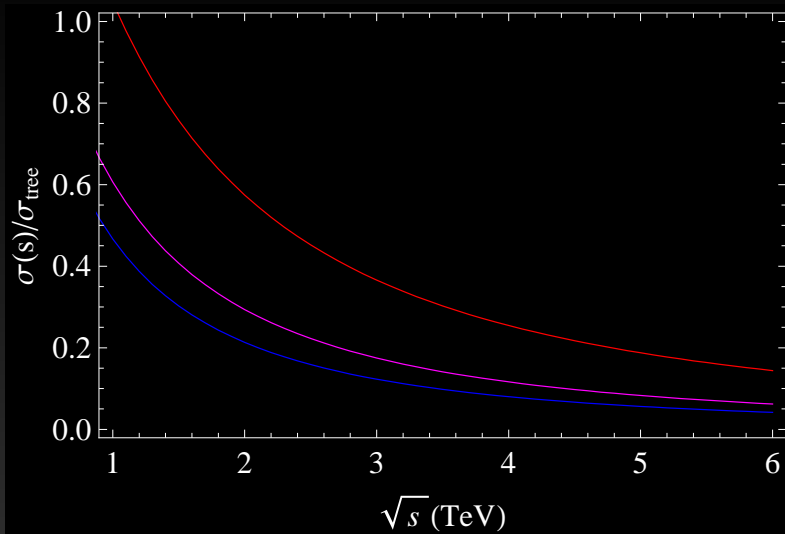


$t = -0.5s$

$t = -0.3s$

$t = -0.1s$

$$u\bar{u}_L \rightarrow t\bar{t}_L$$

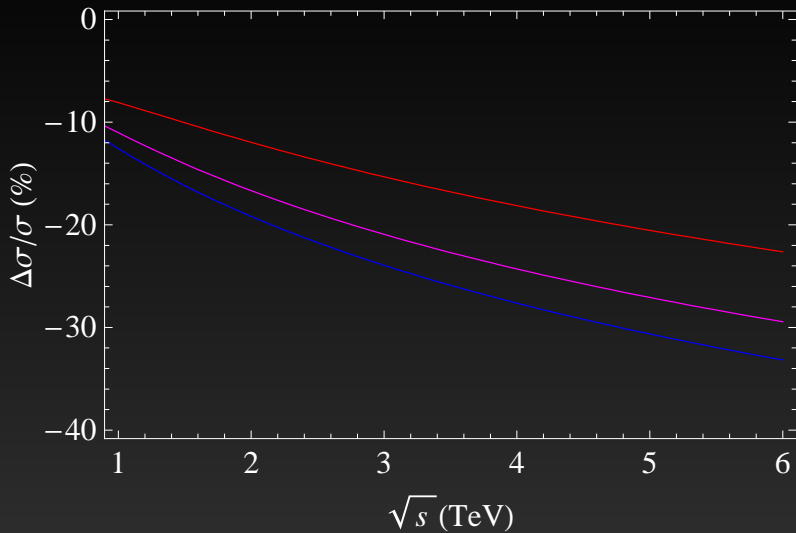


$t=-0.5s$

$t=-0.3s$

$t=-0.1s$

# $u\bar{u} \rightarrow t\bar{t}_L$ , contribution from EW

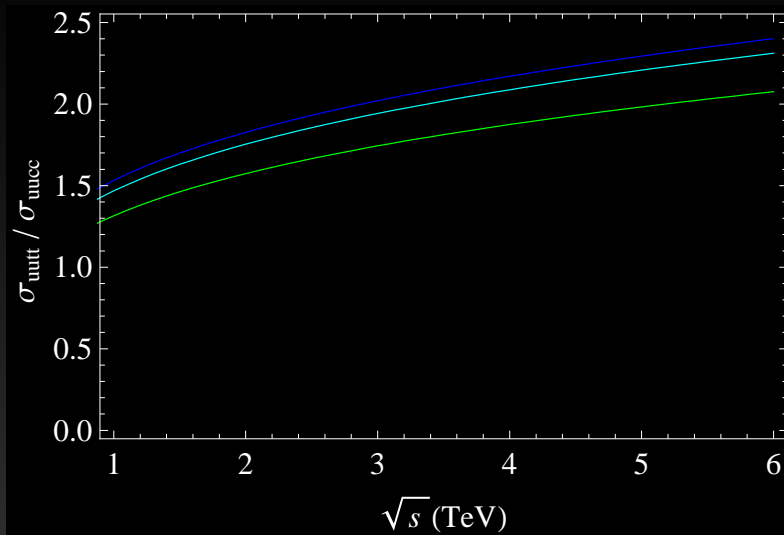


$t = -0.5s$

$t = -0.3s$

$t = -0.1s$

# heavy quark mass effect



$t = -0.5s$

$t = -0.3s$

$t = -0.1s$

# Conclusion

- Processes calculated in toy model:
  - ▶  $u\bar{u} \rightarrow \tilde{q}\tilde{q}^*$
  - ▶  $u\bar{u} \rightarrow c\bar{c}$
- Processes calculated in SM using SCET including
  - ▶ Sudakov form factor for electrons, u and t quarks
  - ▶  $u\bar{u} \rightarrow c\bar{c}$
  - ▶  $u\bar{u} \rightarrow t\bar{t}$
- Ability to include mass effects due to differences in  $M_Z$ ,  $M_W$ ,  $M_h$  and  $m_t$ .
- The results have been easily generalized to include:
  - ▶ production of SUSY particles
  - ▶ 4-particle processes in ILC ( $e^+e^- \rightarrow u\bar{u}, t\bar{t}, \tilde{t}\tilde{t}^*$  etc)

# Standard Model (heavy quarks):

